# Metal - Mott insulator heterostructures: A dynamical mean field study

Karol Makuch, Krzysztof Byczuk



Faculty of Physics, Institute of Theoretical Physics, University of Warsaw, Poland

## Motivation and aims

Phenomena observed at the interface of two materials can be different than those seen in the bulk. For example, in oxide heterostructures two-dimensional electron gases can appear at the interfaces of band and Mott insulator.



## Method: dynamical mean field theory

• W. Metzner, D. Vollhardt, Phys. Rev. Lett. 59, 121 (1987) • A. Georges, G. Kotliar, W. Krauth, M. J. Rozenberg, Rev. Mod. Phys. 68, 13 (1996) •Freericks, J. K. Transport in multilayered nanostructures, Imperial College Press, 2006

 $\boldsymbol{G}_{\mathbf{ij}}(\tau,\tau') = -\langle c_{\mathbf{i}\sigma}(\tau) \hat{c}^{\dagger}_{\mathbf{i}\sigma}(\tau') \rangle$ 

Dyson equation with use of translational symmetry in x and y direction (quasi 1-dimensional system):

 $\left(i\omega_{n}+\mu_{\alpha}-\varepsilon^{||}-\Sigma_{\alpha}\left(i\omega_{n}\right)\right)G_{\alpha\beta}\left(i\omega_{n},\varepsilon^{||}\right)+t_{\alpha\alpha+1}^{1D}G_{\alpha+1\beta}\left(i\omega_{n},\varepsilon^{||}\right)+t_{\alpha\alpha-1}^{1D}G_{\alpha-1\beta}\left(i\omega_{n},\varepsilon^{||}\right)=\delta_{\alpha\beta}$ 

S. Thiel et al., Science 313, 1942 (2006)

We consider three dimensional heterostructure consisted of Mott insulator (green colour below) between two semi-infinite band insulators



$$G_{\alpha\alpha}(i\omega_n) = \int_{-\infty}^{\infty} d\varepsilon \rho^{2d} \left(\varepsilon^{||}\right) G_{\alpha\alpha}\left(i\omega_n,\varepsilon^{||}\right)$$

$$Local Dyson equation:$$

$$\Delta_{\alpha}(i\omega_n) = i\omega_n + \mu - \frac{1}{G_{\alpha\alpha}(i\omega_n)} - \Sigma_{\sigma}(i\omega_n)$$

$$Density of states for square lattice$$

$$\int_{-\infty}^{0.4} \frac{1}{2} \int_{-\infty}^{0.2} \frac{1}{2} \int_{-\frac{1}{2}}^{0.2} \frac{1}{2}$$

Single impurity problem (solved by continuous time Monte Carlo algorithm):

$$G_{\alpha\alpha}(i\omega_{n}) = -\frac{1}{\Xi} \left[ \prod_{\sigma} \int D\varphi_{\sigma} \int D\varphi_{\sigma}^{*} \right] \varphi_{\sigma}(i\omega_{n}) \varphi_{\sigma}^{*}(i\omega_{n}) \exp\left[-S_{\alpha}\right]$$

$$S_{\alpha} = \sum_{\sigma} \int_{0}^{\beta} d\tau \varphi_{\sigma}^{*}(\tau) \left(\frac{\partial}{\partial \tau} - \mu_{\alpha}\right) \varphi_{\sigma}(\tau) + U_{\alpha} \int_{0}^{\beta} d\tau \varphi_{\uparrow}^{*}(\tau) \varphi_{\uparrow}(\tau) \varphi_{\downarrow}^{*}(\tau) \varphi_{\downarrow}(\tau) + \int_{0}^{\beta} d\tau \int_{0}^{\beta} d\tau' \sum_{\sigma} \varphi_{\sigma}^{*}(\tau) \Delta_{\alpha}(\tau - \tau') \varphi_{\sigma}(\tau')$$

$$\beta = 1/(k_{B}T)$$

### Results

Before heterostructure - homogeneous Hubbard model

 $H = t \quad \sum \quad \sum \hat{c}^{\dagger}_{\mathbf{i}\sigma} \hat{c}_{\mathbf{i}\sigma} + U \sum \hat{n}_{\mathbf{i}\uparrow} \hat{n}_{\mathbf{i}\downarrow} - \mu \sum \hat{n}_{\mathbf{i}\sigma} \hat{n}_{\mathbf{i}\sigma}$ 

Our aim is to calculate electronic properties in the heterostructure. In particular:

-quasi particle density of states

- -average kinetic energy of particles
- -double occupation number

which depend on the position in the heterostructure.

## Model of the heterostructure

The system under consideration is described by grand canonical Hamiltonian of one band Hubbard model:

$$H = \sum_{(\mathbf{i},\mathbf{j})\in n.n.} \sum_{\sigma=\uparrow,\downarrow} t_{\mathbf{i}\mathbf{j}} \hat{c}^{\dagger}_{\mathbf{i}\sigma} \hat{c}_{\mathbf{i}\sigma} + \sum_{\mathbf{i}} U_{\mathbf{i}} \hat{n}_{\mathbf{i}\uparrow} \hat{n}_{\mathbf{i}\downarrow} - \sum_{\mathbf{i}} \sum_{\sigma=\uparrow,\downarrow} \mu_{\mathbf{i}} \hat{n}_{\mathbf{i}\sigma}$$
$$\hat{n}_{i\sigma} \equiv \hat{c}^{\dagger}_{i\sigma} \hat{c}_{i\sigma}$$

for infinite simple cubic lattice

 $n_x, n_y, n_z = \dots, -1, 0, 1, \dots$  $\mathbf{j}, \mathbf{i} \in [n_x \hat{\mathbf{e}}_x, n_y \hat{\mathbf{e}}_y, n_z \hat{\mathbf{e}}_z]$ 

with the following properties  
-uniform hopping in all directions: 
$$t_{ij} = t$$
  
-assumed half-filling  $\mu_i = \frac{U_i}{2}$ 

#### $(\mathbf{i},\mathbf{j}) \in n.n. \sigma = \uparrow,\downarrow$

Density of states (at chemical potential)



$$G(\tau) = \int_{-\infty}^{\infty} d\omega \rho(\omega) \frac{e^{-(\omega-\mu)\tau}}{1 + e^{-(\omega-\mu)\beta}}$$

Average order expansion in Monte Carlo simulation (kinetic energy)

 $\sigma = \uparrow, \downarrow$ 



Kinetic energy by average expansion order in continuous time – hybridization expansion Monte Carlo simulations

 $\epsilon_K = -k_B T < k >$ 

Haule, K., 2007, Phys. Rev. B 75, 155113

Results for heterostructure

10

5

15



-interaction depends on lattice site in the heterostructure (translational symmetry in x and y direction)



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